

MIT OpenCourseWare
<http://ocw.mit.edu>

6.013/ESD.013J Electromagnetics and Applications, Fall 2005

Please use the following citation format:

Markus Zahn, *6.013/ESD.013J Electromagnetics and Applications, Fall 2005*. (Massachusetts Institute of Technology: MIT OpenCourseWare). <http://ocw.mit.edu> (accessed MM DD, YYYY). License: Creative Commons Attribution-Noncommercial-Share Alike.

Note: Please use the actual date you accessed this material in your citation.

For more information about citing these materials or our Terms of Use, visit:
<http://ocw.mit.edu/terms>

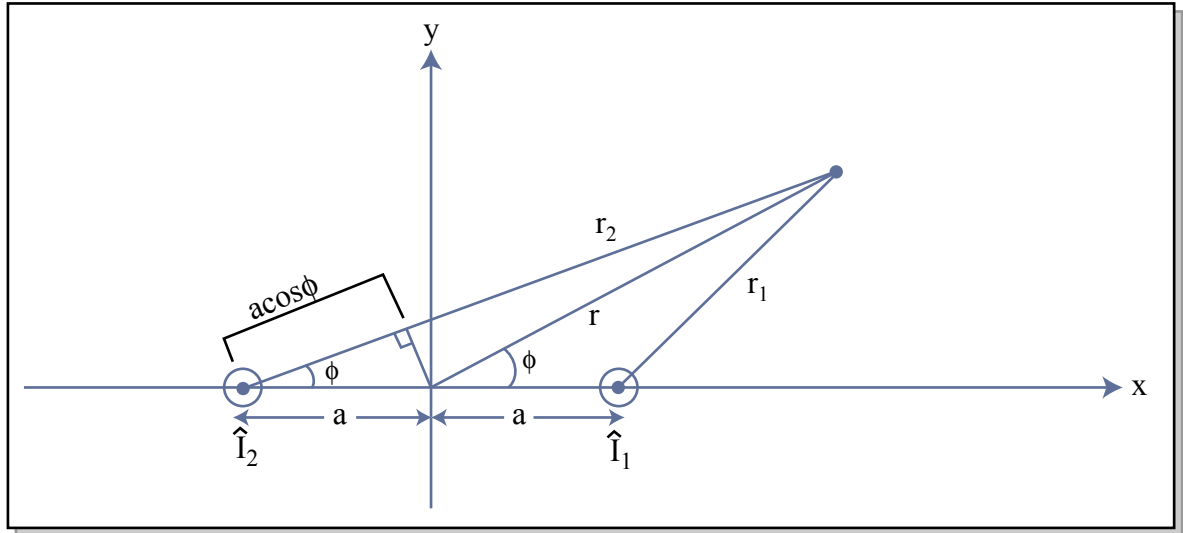
I. Two Element Array in $\theta = \frac{\pi}{2}$ plane (x-y plane)

Image by MIT OpenCourseWare.

Far field ($kr \gg 1, r \gg a$)

$$\hat{E}_\theta(r, \theta = \frac{\pi}{2}, \phi) = \frac{\hat{E}_1}{jkr_1} e^{-jkr_1} + \frac{\hat{E}_2}{jkr_2} e^{-jkr_2} = \eta \hat{H}_\phi(r, \theta = \frac{\pi}{2}, \phi)$$

$$\hat{E}_1 = -\frac{\hat{I}_1 dl k^2 \eta}{4\pi}$$

$$\hat{E}_2 = -\frac{\hat{I}_2 dl k^2 \eta}{4\pi}$$

$$r_2 \approx r + a \cos(\phi), r_1 \approx r - a \cos(\phi)$$

$$\hat{E}_\theta(r, \theta = \frac{\pi}{2}, \phi) = \eta \hat{H}_\phi(r, \theta = \frac{\pi}{2}, \phi) \approx \underbrace{-\frac{k^2 \eta dl}{4\pi jkr}}_{\text{element factor}} e^{-jkr} \left[\hat{I}_1 e^{+jka \cos(\phi)} + \hat{I}_2 e^{-jka \cos(\phi)} \right]_{\text{array factor}}$$

Assume: $\hat{I}_1 = \hat{I}, \hat{I}_2 = \hat{I} e^{j\chi} \Rightarrow \hat{E}_1 = \hat{E}_0, \hat{E}_2 = \hat{E}_0 e^{j\chi}$

$$\begin{aligned} E_\theta(r, \theta = \frac{\pi}{2}, \phi) &= \eta \hat{H}_\phi(r, \theta = \frac{\pi}{2}, \phi) = \frac{\hat{E}_0}{jkr} e^{-jkr} \left[e^{+jka \cos(\phi)} + e^{j\chi} e^{-jka \cos(\phi)} \right] \\ &= \frac{\hat{E}_0}{jkr} e^{-jkr} e^{j\chi/2} \left[\underbrace{e^{-j(\frac{\chi}{2} - ka \cos(\phi))} + e^{j(\frac{\chi}{2} - ka \cos(\phi))}}_{2 \cos(-\frac{\chi}{2} + ka \cos(\phi))} \right] \\ &= \frac{2\hat{E}_0}{jkr} e^{-jkr} e^{j\chi/2} \cos\left(-\frac{\chi}{2} + ka \cos(\phi)\right) \end{aligned}$$

$$\langle S_r(t, \theta = \frac{\pi}{2}, \phi) \rangle = \frac{1}{2} \frac{|\hat{E}_\theta|^2}{\eta} = \frac{2|\hat{E}_0|^2}{\eta(kr)^2} \cos^2\left(ka \cos(\phi) - \frac{\chi}{2}\right)$$

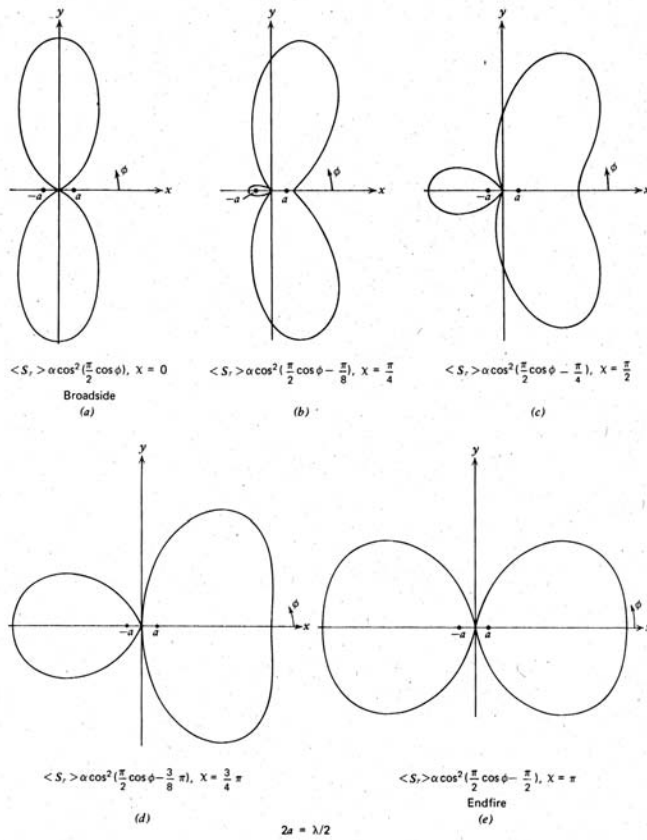


Figure 9-7 The power radiation pattern due to two-point dipoles depends strongly on the dipole spacing and current phases. With a half wavelength dipole spacing ($2a = \lambda/2$), the radiation pattern is drawn for various values of current phase difference in the $\theta = \pi/2$ plane. The broadside array in (a) with the currents in phase ($\chi = 0$) has the power lobe in the direction perpendicular to the array while the end-fire array in (e) has out-of-phase currents ($\chi = \pi$) with the power lobe in the direction along the array.

Broadside:

$$2a = \frac{\lambda}{2}, \chi = 0, ka = \frac{\pi}{2}$$

$$\langle S_r \rangle = \frac{2|\hat{E}_0|^2}{\eta(kr)^2} \cos^2\left(\frac{\pi}{2} \cos(\phi)\right)$$

Endfire:

$$2a = \frac{\lambda}{2}, \chi = \pi, ka = \frac{\pi}{2}$$

$$\langle S_r \rangle = \frac{2|\hat{E}_0|^2}{\eta(kr)^2} \cos^2\left(\frac{\pi}{2} (\cos(\phi) - 1)\right)$$

From *Electromagnetic Field Theory: A Problem Solving Approach*, by Markus Zahn, 1987. Used with permission.

$$\text{Maxima: } ka \cos(\phi) - \frac{\chi}{2} = \pm m\pi, m = 0, 1, 2, \dots$$

$$\text{Minima: } ka \cos(\phi) - \frac{\chi}{2} = \pm(2m + 1)\frac{\pi}{2}, m = 0, 1, 2, \dots$$

Case Studies:

$$2a = \frac{\lambda}{2} \Rightarrow ka = \frac{2\pi a}{\lambda} = \frac{2\pi a}{4a} = \frac{\pi}{2}$$

$$2a = \lambda \Rightarrow ka = \pi$$

$$2a = \frac{\lambda}{2} \Rightarrow \frac{\pi}{2} \cos(\phi) - \frac{\chi}{2} = \pm m\pi \text{ (maxima)} \Rightarrow \cos(\phi) = \frac{\chi}{\pi} \pm 2m$$

$$\frac{\pi}{2} \cos(\phi) - \frac{\chi}{2} = \pm(2m + 1)\frac{\pi}{2} \text{ (minima)} \Rightarrow \cos(\phi) = \frac{\chi}{\pi} \pm (2m + 1)$$

$$\lambda = 4a$$

| χ | $\cos(\phi_{\max})$ | $\cos(\phi_{\min})$ | ϕ_{\max} | ϕ_{\min} | |
|------------------|---------------------|---------------------|--------------------|-------------------|-----------|
| 0 | 0 | 1 | $\pm\frac{\pi}{2}$ | $0, \pi$ | Broadside |
| $\frac{\pi}{4}$ | $\frac{1}{4}$ | $-\frac{3}{4}$ | $\pm 75.5^\circ$ | $\pm 138.6^\circ$ | |
| $\frac{\pi}{2}$ | $\frac{1}{2}$ | $-\frac{1}{2}$ | $\pm 60^\circ$ | $\pm 120^\circ$ | |
| $\frac{3\pi}{4}$ | $\frac{3}{4}$ | $-\frac{1}{4}$ | $\pm 41.4^\circ$ | $\pm 104.5^\circ$ | |
| π | 1 | 0 | $0, \pi$ | $\pm 90.^\circ$ | Endfire |

$$\lambda = 2a \Rightarrow \cos(\phi_{\max}) = \frac{\chi}{2\pi} \pm m, \cos(\phi_{\min}) = \frac{\chi}{2\pi} \pm \frac{1}{2}(2m + 1)$$

| χ | $\cos(\phi_{\max})$ | $\cos(\phi_{\min})$ | ϕ_{\max} | ϕ_{\min} |
|------------------|-----------------------------|-----------------------------|------------------------------|---------------------------|
| 0 | 0, 1 | $\frac{1}{2}, -\frac{1}{2}$ | $0, \pm 90^\circ, 180^\circ$ | $\pm 60^\circ$ |
| $\frac{\pi}{4}$ | $\frac{1}{8}, -\frac{7}{8}$ | $-\frac{3}{8}, \frac{5}{8}$ | $82.8^\circ, 151^\circ$ | $51^\circ, 112^\circ$ |
| $\frac{\pi}{2}$ | $\frac{1}{4}, -\frac{3}{4}$ | $-\frac{1}{4}, \frac{3}{4}$ | $75.5^\circ, 138.6^\circ$ | $41.4^\circ, 104.5^\circ$ |
| $\frac{3\pi}{4}$ | $\frac{3}{8}, -\frac{5}{8}$ | $-\frac{1}{8}, \frac{7}{8}$ | $68.0^\circ, 128.7^\circ$ | $29.0^\circ, 97.2^\circ$ |
| π | $\frac{1}{2}, -\frac{1}{2}$ | 0, 1 | $60^\circ, 120^\circ$ | $90^\circ, 0^\circ$ |

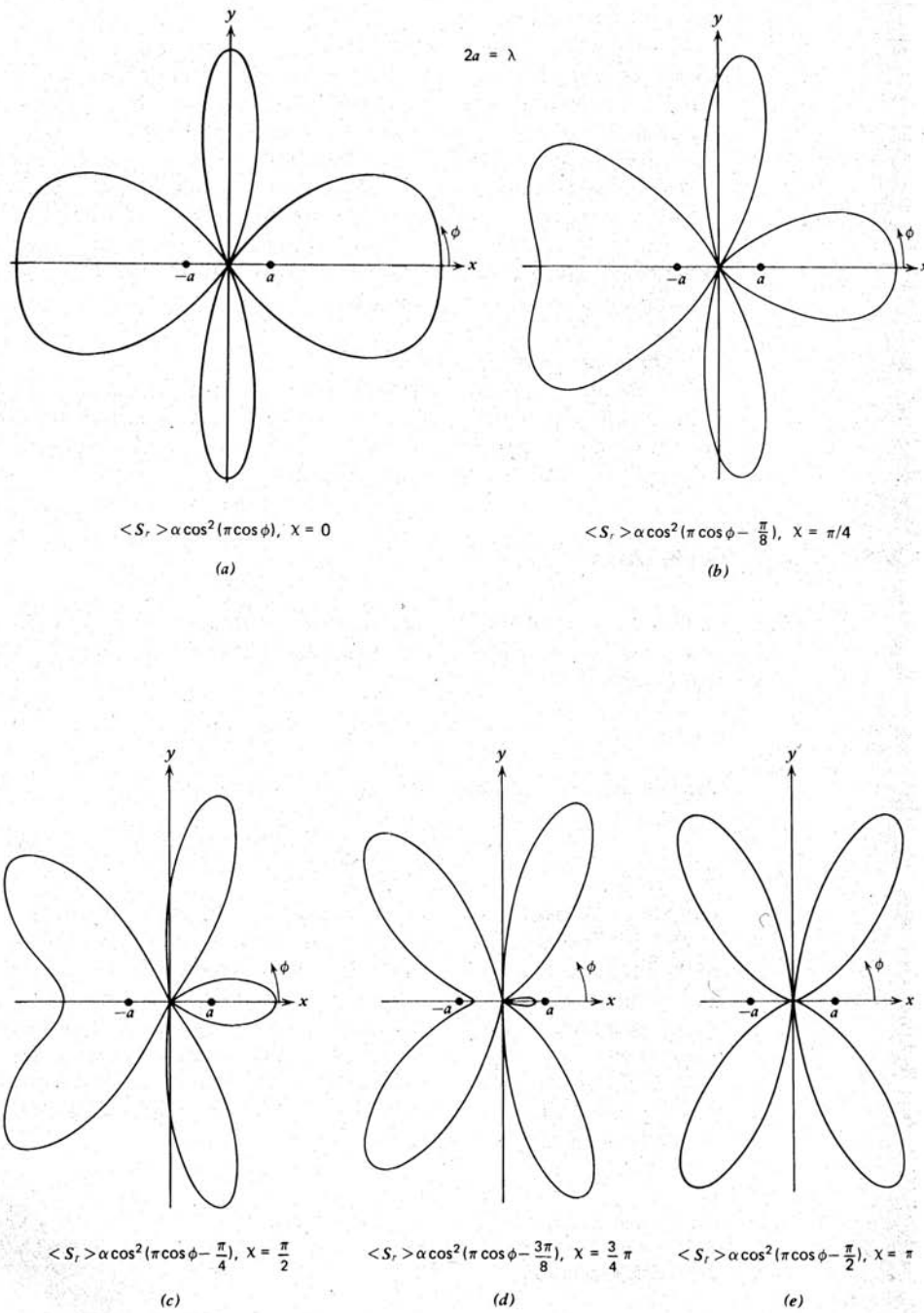


Figure 9-8 With a full wavelength dipole spacing ($2a = \lambda$) there are four main power lobes.

From *Electromagnetic Field Theory: A Problem Solving Approach*, by Markus Zahn, 1987. Used with permission.

II. An N Dipole Array ($\theta = \frac{\pi}{2}$)

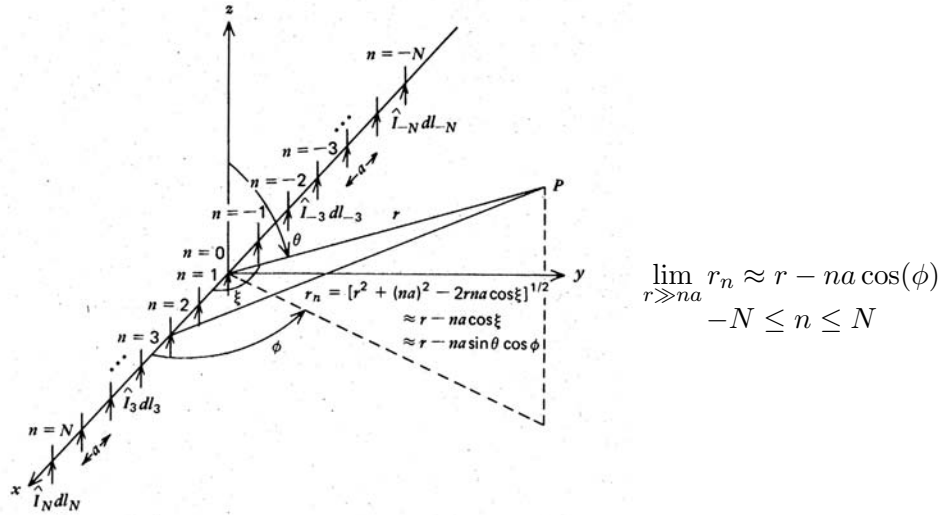


Figure 9-9 A linear point dipole array with $2N+1$ equally spaced dipoles.

From *Electromagnetic Field Theory: A Problem Solving Approach*, by Markus Zahn, 1987. Used with permission.

$$\begin{aligned}\hat{E}_\theta \left(r, \theta = \frac{\pi}{2}, \phi \right) &= \eta \hat{H}_\phi \left(r, \theta = \frac{\pi}{2}, \phi \right) \\ &= -\frac{k\eta dl}{4\pi jr} \underbrace{\left[\sum_{-N}^{+N} \hat{I}_n e^{jkna \cos(\phi)} \right]}_{\text{Array factor} = \text{AF}} e^{-jkr}\end{aligned}$$

Example:

$$\hat{I}_n = I_0 e^{-jn\chi_0} \quad -N \leq n \leq N$$

$$\text{AF} = I_0 \sum_{-N}^{+N} e^{jn(ka \cos(\phi) - \chi_0)}$$

$$\text{Let } \beta \equiv e^{j(ka \cos(\phi) - \chi_0)}$$

$$S = \frac{\text{AF}}{I_0} = \sum_{-N}^{+N} \beta^n = \beta^{-N} + \beta^{-N+1} + \dots + \beta^{-2} + \beta^{-1} + 1 + \beta + \beta^2 + \dots + \beta^{N-1} + \beta^N$$

$$S(1 - \beta) = \beta^{-N} - \beta^{N+1} \Rightarrow S = \frac{\beta^{-N} - \beta^{N+1}}{1 - \beta} = \frac{\beta^{-N-\frac{1}{2}} - \beta^{N+\frac{1}{2}}}{\beta^{-1/2} - \beta^{1/2}}$$

multiply by $\frac{\beta^{-1/2}}{\beta^{-1/2}}$

$$S = \frac{\sin \left[\left(N + \frac{1}{2} \right) (ka \cos(\phi) - \chi_0) \right]}{\sin \left[\frac{1}{2} (ka \cos(\phi) - \chi_0) \right]}$$

Maxima: $ka \cos(\phi) - \chi_0 = 2n\pi, n = 0, 1, 2, \dots$

Principal maximum at $n = 0 \Rightarrow \cos(\phi) = \frac{\chi_0}{ka}$

Minima: $(N + \frac{1}{2})(ka \cos(\phi) - \chi_0) = n\pi, n = 1, 2, 3, \dots$

Demonstration, $N = 2$ (2 dipole array)

$$2a = \frac{3}{2}\lambda, \chi_0 = 0$$

$$I \propto \cos^2(ka \cos(\phi)) = \cos^2\left(\frac{2\pi}{\lambda} \frac{3}{2}\lambda \cos(\phi)\right) = \cos^2\left(\frac{3\pi}{2} \cos(\phi)\right)$$

$$\text{Minima: } \frac{3\pi}{2} \cos(\phi) = \frac{\pi}{2} \Rightarrow \cos(\phi) = \frac{1}{3} \Rightarrow \phi = 70.5^\circ$$

$$\frac{3\pi}{2} \cos(\phi) = \frac{3\pi}{2} \Rightarrow \cos(\phi) = 1 \Rightarrow \phi = 0^\circ$$

$$\text{Maxima: } \frac{3\pi}{2} \cos(\phi) = 0 \Rightarrow \phi = 90^\circ$$

$$\frac{3\pi}{2} \cos(\phi) = \pi \Rightarrow \cos(\phi) = \frac{2}{3} \Rightarrow \phi = 48.2^\circ$$

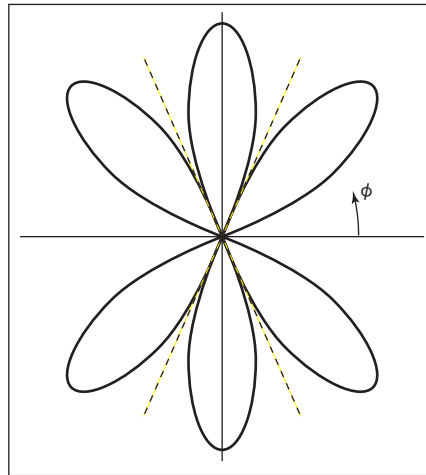


Image by MIT OpenCourseWare.

Intensity pattern

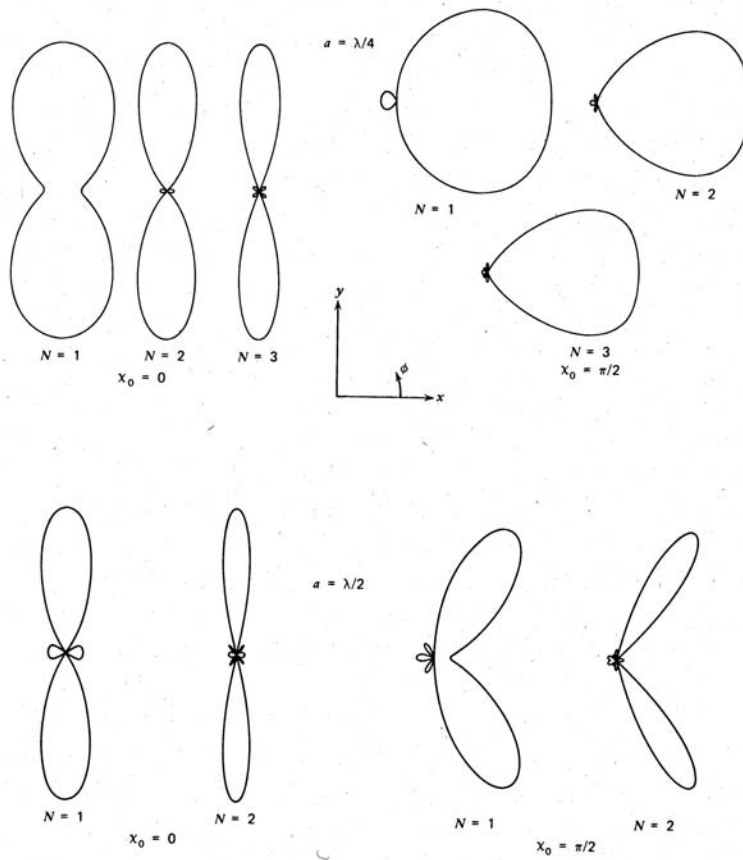


Figure 9-10 The radiation pattern for an N dipole linear array for various values of N , dipole spacing $2a$, and relative current phase χ_0 in the $\theta = \pi/2$ plane.

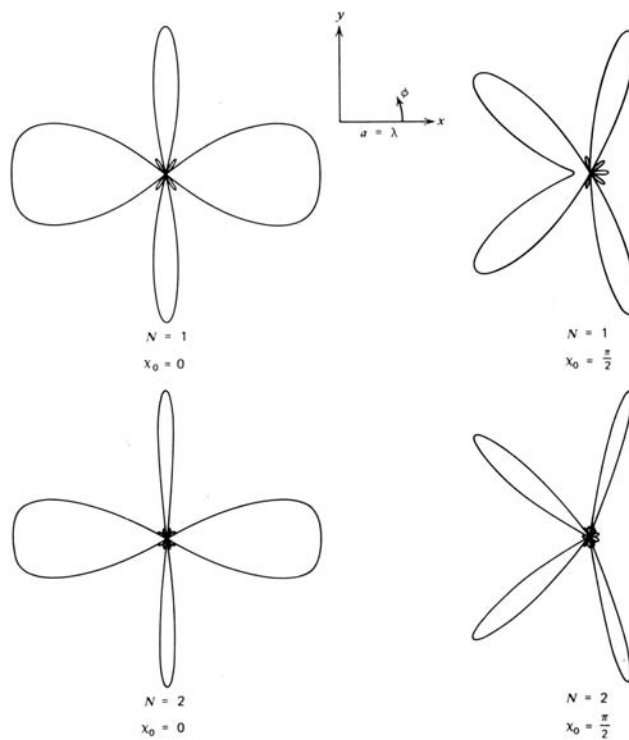


Figure 9-10

From *Electromagnetic Field Theory: A Problem Solving Approach*, by Markus Zahn, 1987. Used with permission.